# DEVICE-INDEPENDENT QUANTUM KEY DISTRIBUTION WITH SINGLE-PHOTON SOURCES

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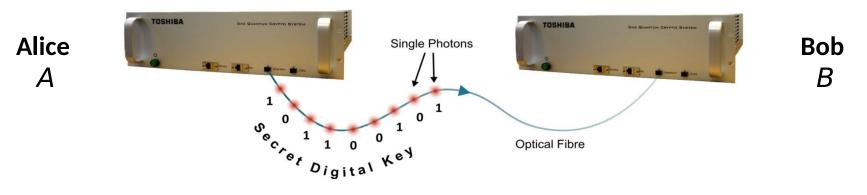
## WHAT IS THE CURRENT STATUS OF QUANTUM KEY DISTRIBUTION?

#### IT IS COMMERCIALLY AVAILABLE...

COMPANIES SUCH AS TOSHIBA, MAGIQ, ID QUANTIQUE... :



Secret Digital Key Exchange Using Quantum Key Distribution



Commercial devices implementing BB84, SARG, COW (iDQ),... protocols.

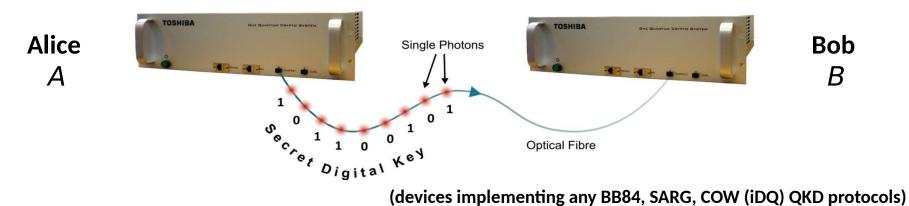
#### **OK, BUT THE ABOVE IMPLEMENTATIONS HAVE BEEN**

"HACKED"? RESEARCHERS WERE ABLE TO EAVESDROP AND CAPTURE THE KEY WITHOUT LEAVING IT IS THE <u>DEVICES</u> THAT HAVE BEEN CRACKED AND NOT THE CONCEPT OF



# **"HACKING" QUANTUM KEY DISTRIBUTION**

#### Secret Digital Key Exchange Using Quantum Key Distribution

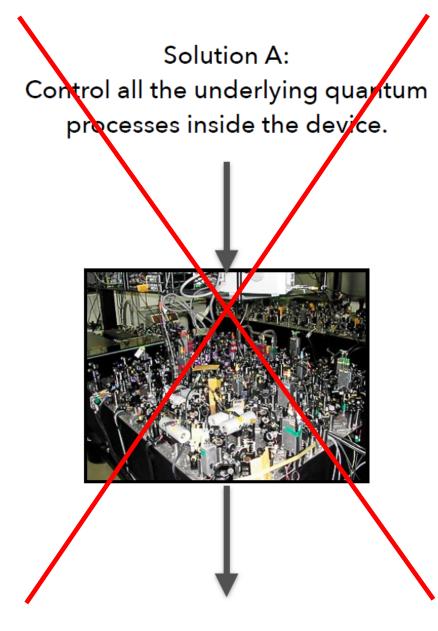


- After the end of the key distribution protocol devices announce:
  "A secure key has been successfully established and reads #\$%#\$&...."
- This means that the **error rate** has been verified to be below a certain threshold ( $< \epsilon$ %), which "guarantees" by laws of quantum physics that no-one can have access to the key.
- Ok, but this ε% is derived assuming a particular (quantum mechanical) model of the devices importantly modelling: optical fibres, detectors, electronics, losses, detection inefficiencies etc.
- HACKING:

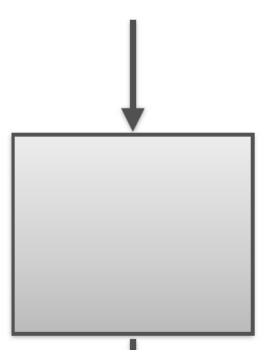
Explore other <u>degrees of freedom</u> that are not accounted for in the model, whose presence invalidates the proofs of security.

#### IT IS ALL ABOUT THE MISMATCH BETWEEN THE THEORETICAL REQUIREMENTS AND THE IMPLEMENTATION!

# **IS THERE A WAY AROUND THIS?**



Solution B: Make no assumptions about the internal working of the device.



# SOLUTION B: DEVICE-INDEPENDENT (DI) APPROACH

Treat the devices as **black boxes** with:

- input buttons {*x*,*y*} (QKD: randomly chosen measurement settings)

- **output** bulbs {*a*,*b*} (QKD: outcomes of the implemented measurements):

Assure the security basing on the **probability distribution (***behaviour***)** P(a, b|x, y) that Alice and Bob may reconstruct from some subset of data using the classical authorised channel (*they call one another*).

This is possible as P(a, b|x, y) ideally exhibits **non-local correlations** that cannot be explained with **classical physics** but only with **quantum mechanics**  $\rightarrow$  **Bell violation**.

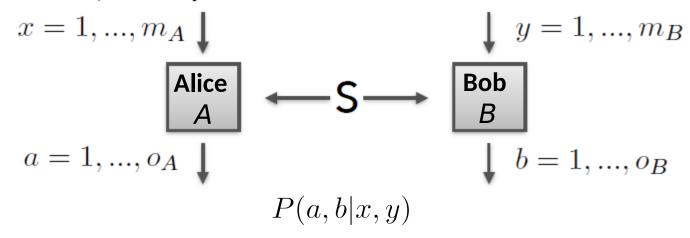
#### **DRAWBACK:**

Such approach is very **sensitive to noise**. After introducing imperfections (*transmission*, *detection losses*, *etc.*) in devices, the correlations quickly become classically explainable (detection loophole).

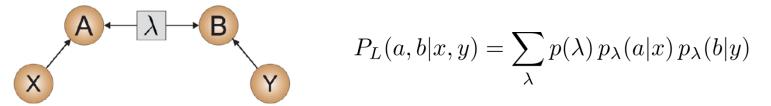
[Acin et al. Phys. Rev. Lett. 97 (2007), Barrett et al., Phys. Rev. Lett. 95 (2005)]

# **BELL VIOLATION IN 2 SLIDES**

At each round of the test, Alice and Bob perform measurements x and y on some part of a system S and retrieve outcomes a and b:



• **Classical** explanation of correlations – **Local Hidden Variable Model** (LHVM):

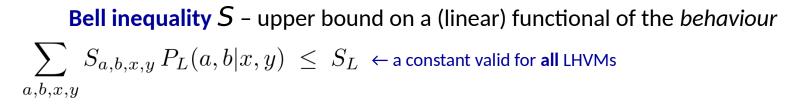


• Quantum mechanics allows for stronger nonlocal correlations to be shared:

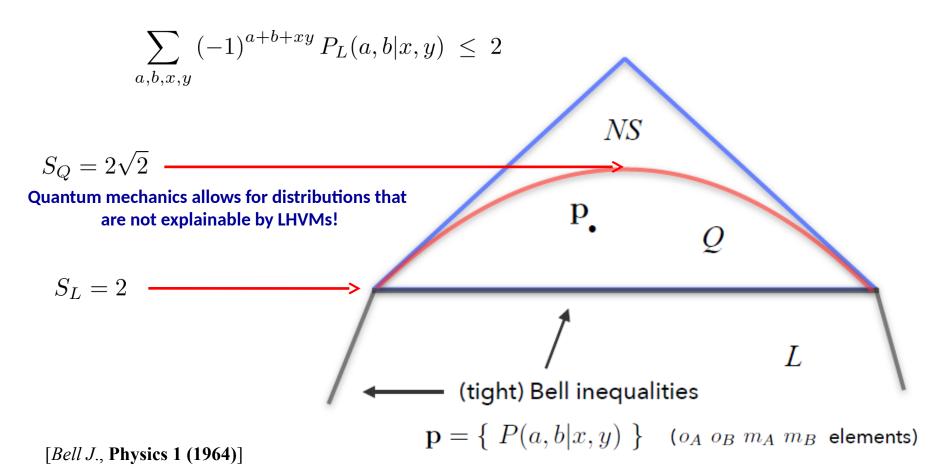
$$P_Q(a,b|x,y) = \operatorname{Tr} \left\{ \rho_{AB} \ M_{a|x} \otimes M_{b|y} \right\}$$

(...much richer structure)

# **BELL INEQUALITIES - GEOMETRIC REPRESENTATION**



e.g., CHSH inequality (2 inputs, 2 outcomes):  $S_{a,b,x,y} = (-1)^{a+b+xy}$   $S_L = 2$ 



### **GUESSING PROBABILITY OF AN EAVESDROPPER**

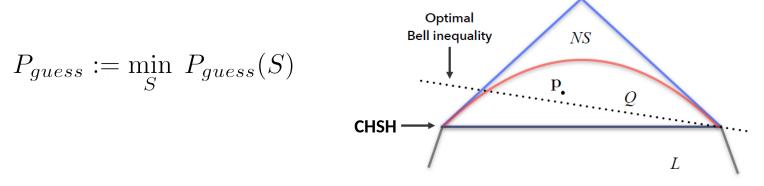
For a given Bell inequality S and its violation  $S_{obs}$  by the observed behaviour  $\mathbf{p}_{obs}$ 

one can explicitly calculate the **guessing probability**, i.e.,

The maximal probability that an eavesdropper correctly guesses the outcome of a box (A) can upper bounded for a particular S by:

$$P_{guess}(S) = \max P(a|x)$$
  
s.t. 
$$\begin{cases} \sum_{a,b,x,y} S_{a,b,x,y} P(a,b|x,y) = S_{obs}[\mathbf{p}_{obs}] \\ P(a,b|x,y) \in Q \\ P(a,b|x,y) \in Q \\ \text{Quantum set } Q \text{ is a convex space but not a simplex need Semi-Definite Programming (SDP) tricks, i.e., the NPA hierarchy [Navascues et al., PRL 98 (2007)] \end{cases}$$

Furthermore, one should **optimise over all Bell inequalities** to make the guessing probability (and, hence, the *power of eavesdropper*) as **small** as possible.

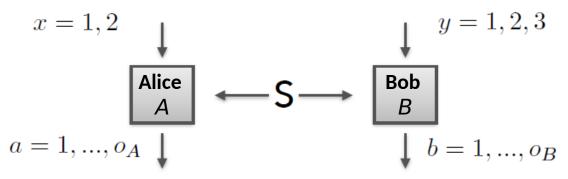


A convex problem – again efficiently solvable by an SDP

[Colbeck R., PhD Thesis, Cambridge (2009); Pironio et al., Nature 464 (2010)]

### KEY RATE IN DEVICE-INDEPENDENT QUANTUM KEY DISTRIBUTION

#### **DI-QKD protocol:**

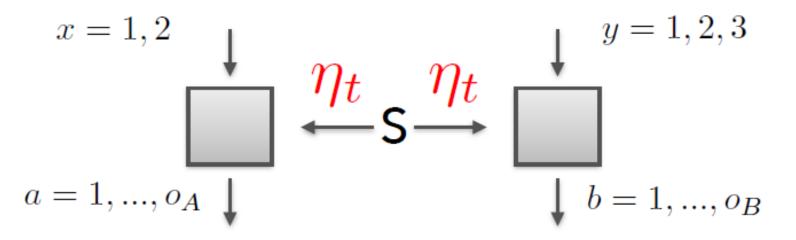


- All rounds for which  $y \neq 3$  are used to generate  $\mathbf{p} = \{P(a, b | x, y)\}$ .
- From  $\mathbf{p}$  Alice and Bob construct (as discussed before)  $P_{guess}$ .
- Rounds in which x=1 and y=3 are used to generate the key.
- The key rate of the DI-QKD protocol is lower-bounded by:

 $r \geq -\log_2(P_{guess}) - H(x = 1 | y = 3)$   $r \geq -\log_2(P_{guess}) - H(x = 1 | y = 3)$   $r \geq 0$   $r \geq 0$  r

[*Masanes et al.* Nat. Comms 2 (2011), *Pironio* et al. PRX 3 (2013)] [*Vazirani & Viddick* PRL 113 (2014), *Arnon-Friedman* et al. arXiv:1607.01797]

### **PROBLEM 1: OF TRANSMISSION LOSSES**



Loss in optic fibres decays exponentially with distance:  $\eta_t = e^{-L/L_{att}}$ With  $L_{att} = 22 \ km$ , for a distance  $L = 10 \ km$  we have:  $\eta_t \approx 60\%$ 

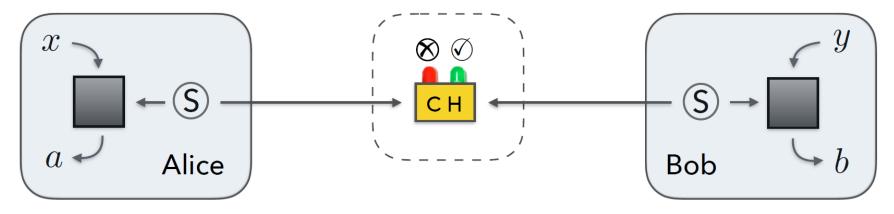
- 66% is the fundamental limit to violate any Bell Inequality (not to mention DIQKD)
- For long distance communications we should not hope for technological progress to resolve the problem...

# SOLUTION 1: HERALDING (WITHOUT OPENING THE DETECTION LOOPHOLE)

Side-Heralding (a'la amplification):



**Central-Heralding** (a'la entanglement swapping, quantum repeaters):



#### **PROBLEM 2:**

#### "STANDARD HERALDING" WITH SPONTANEOUS PHOTON-PAIR SOURCES IS OF NO USE!

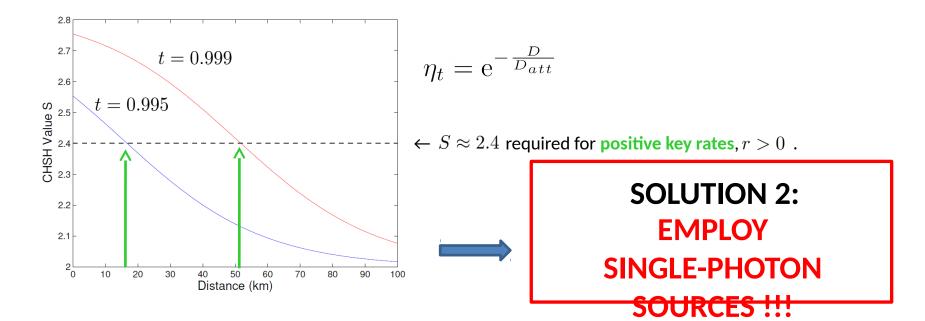
Imagine that Alice and Bob share (inside the boxes) an **entangled** photon-pair produced in **spontaneous parametric down-conversion (SPDC)** process with heralding implemented via, e.g., "qubit amplification" ( t) [Gikin et al. PRL 105 (2010)]:

$$\rho_{AB} = (1-p)(1-t)^2 |vac\rangle \langle vac| + p \eta_t t(1-t) |\psi_-\rangle \langle \psi_-|$$

 $S = 2 \left[ 1 + \frac{p \eta_t t (\sqrt{2} - 1)}{(1 - p)(1 - t) + p \eta_t t} \right]$ 

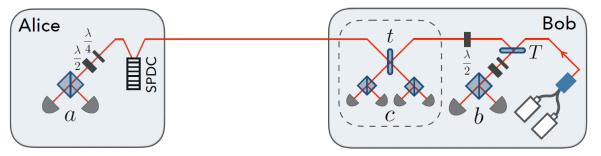
Let us consider the CHSH value of Bell violation:

CHSH value *S* exponentially quickly approaches the local value  $S_L = 2$  with distance. [intuition: vacuum terms always eventually dominate as they are O(1) in p.]

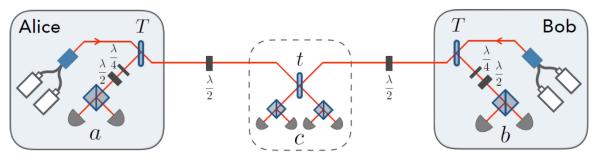


### **DIQKD SCHEMES WITH SIDE- AND CENTRAL-HERALDING**

Side-Heralding (1 SPDC, 2 SPSs) ["Qubit amplifier" inspired by Pitkanen et al. PRA 84 (2011)]:



Central-Heralding (4 SPSs) ["Quantum repeater" inspired by Lasota et al. PRA 90 (2014)]:



**OPTIMIZING PARAMETERS** t, T FOR EACH SCHEME, ASSUMING SOURCES:  $|n\rangle\langle n|$  (with  $p = 10^{-4}$ )

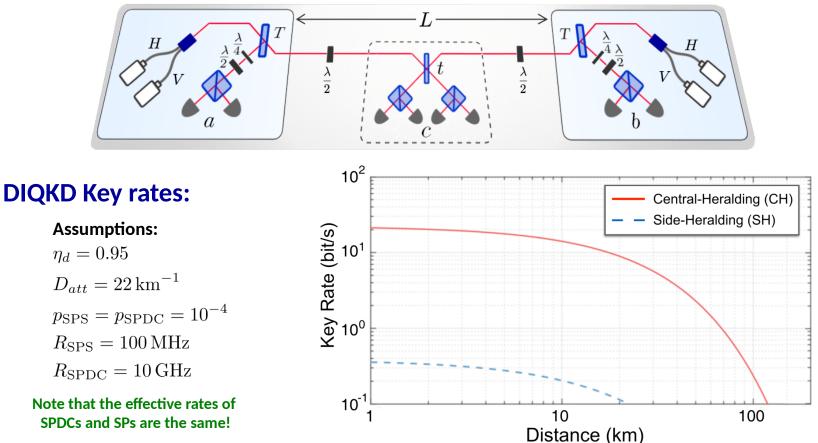
n=1

DIQKD Scheme:	Side-heralding (SH)	Central-heralding (CH)
Critical detection efficiency $\eta_{\rm d}^{\star}$ (diqkd) Critical detection efficiency $\eta_{\rm d}^{\star}$ (nonloc.)	$94.9\%\74.3\%$	$94.3\% \\ 69.2\%$
Noise robustness (nonloc.) Secret key per heralded round (bit fraction $< 1$ )	$31.2\% \ 0.82$	$35.7\% \\ 0.95$

 $\eta_d \leftarrow$  detection efficiency inside one (Alice or Bob) lab (fibre coupling, transmission to detectors, detectors inefficiencies).

### **DIQKD CH-SCHEME PERFORMANCE**

#### **Central-heralding (CH) scheme:**



### **MAIN MESSAGE:**

If already now we were able to get the effective detection efficiency of a device (i.e., all single-photon creation, source-detector transmission and detection efficiencies combined) up to 95%, we would be able to do DIQKD over 50kms at a rate 1bit/sec!